

# A calculus on Lévy exponents and selfdecomposability on Banach spaces\*

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**ABSTRACT.** In infinite dimensional Banach spaces there is no complete characterization of the Lévy exponents of infinitely divisible probability measures. Here we propose *a calculus on Lévy exponents* that is derived from some random integrals. As a consequence we prove that *each* selfdecomposable measure can be factorized as another selfdecomposable measure and its background driving measure that is s-selfdecomposable. This complements a result from the paper of Iksanov-Jurek-Schreiber in the Annals of Probability **32**, 2004.

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Abbreviated title: *A calculus on Lévy exponents*

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**1. Introduction.** Recall that a Borel probability measures  $\mu$ , on a real separable Banach space  $E$ , is called *infinitely divisible* if for each natural number  $n$  there exists a probability measure  $\mu_n$  such that  $\mu_n^{*n} = \mu$ ; the class of all infinitely divisible measures will be denoted by  $ID$ . It is well-know that their Fourier transforms (*the Lévy-Khintchine formulas*) can be written as follows

$$\begin{aligned} \hat{\mu}(y) &= e^{\Phi(y)}, \quad y \in E', \quad \text{and the exponents } \Phi \text{ are of the form} \\ \Phi(y) &= i \langle y, a \rangle - \frac{1}{2} \langle y, Ry \rangle + \int_{E \setminus \{0\}} [e^{i \langle y, x \rangle} - 1 - i \langle y, x \rangle 1_B(x)] M(dx), \end{aligned} \tag{1}$$

where  $E'$  denote the dual Banach space,  $\langle \cdot, \cdot \rangle$  is an appropriate bilinear form between  $E'$  and  $E$ ,  $a$  is a *shift vector*,  $R$  is a *covariance operator* corresponding to the Gaussian part of  $\mu$  and  $M$  is a *Lévy spectral measure*. There is a one-to-one corresponds between  $\mu \in ID$  and the triples  $[a, R, M]$  in its Lévy-Khintchine formula (1); cf. Araujo-Giné (1980), Chapter 3, Section 6, p. 136. The function  $\Phi(y)$  from (1) is called the *Lévy exponent* of  $\mu$ .

*REMARK 1.* (a) If  $E$  is a Hilbert space then Lévy spectral measures  $M$  are completely characterized by the integrability condition  $\int_E (1 \wedge \|x\|^2) M(dx) < \infty$  and Gaussian covariance operators  $R$  coincide with the positive trace-class operators ; cf. Parthasarathy (1967), Chapter VI, Theorem 4.10.

(b) When  $E$  is an *Euclidean space* then Lévy exponents are completely characterized as continuous negative-definite functions; cf. Cuppens (1975) and Schoenberg's Theorem on p. 80.

Finally, a *Lévy process*  $Y(t), t \geq 0$ , means a continuous in probability process with stationary and independent increments and  $Y(0) = 0$ . Without loss of generality we may and do assume that it has paths in the Skorochod space  $D_E[0, \infty)$  of  $E$ -valued *cadlag functions* (i.e., right continuous with left hand limits). There is a one-to-one correspondence between the class  $ID$  and the class of Lévy processes.

The cadlag paths of a process  $Y$  allows us define *random integrals* of the form  $\int_{(a,b]} h(s) Y(r(ds))$  via the formal formula of integration by parts.

Namely,

$$\int_{(a,b]} h(s)Y(r(ds)) := h(b)Y(r(b)) - h(a)Y(r(a)) - \int_{(a,b]} Y(r(s))dh(s), \quad (2)$$

where  $h$  is a real valued function of bounded variation and  $r$  is a monotone and right-continuous function. Furthermore, we have

$$\mathcal{L}\left(\widehat{\int_{(a,b]} h(s)Y(r(ds))}\right)(y) = \exp \int_{(a,b]} \log \widehat{\mathcal{L}(Y(1))}(h(s)y)dr(s), \quad (3)$$

where  $\mathcal{L}(\cdot)$  denotes the probability distribution and  $\hat{\mu}(\cdot)$  denotes the Fourier transform of a measure  $\mu$ ; cf. Jurek-Vervaat (1983) or Jurek (1985) or Jurek-Mason (1993), Section 3.6, p. 116.

**2. A calculus on Lévy exponents.** Let  $\mathcal{E}$  denotes the totality of all functions  $\Phi : E' \rightarrow \mathbb{C}$  appearing as the exponent in the Lévy-Khintchine formula (1). Hence we have that

$$\mathcal{E} + \mathcal{E} \subset \mathcal{E}, \quad \lambda \cdot \mathcal{E} \subset \mathcal{E}, \quad \text{for all postive } \lambda, \quad (4)$$

which means that  $\mathcal{E}$  forms a cone in the space of all complex valued functions defined on  $E'$ . Furthermore, if  $\Phi \in \mathcal{E}$  then all dilations  $\Phi(a\cdot) \in \mathcal{E}$ . These follow from the fact that infinite divisibility is preserved under convolution and under (convolution) powers to positive real numbers.

Here we consider two integral operators acting on  $\mathcal{E}$  or its part. Namely,

$$\begin{aligned} \mathcal{J} : \mathcal{E} &\rightarrow \mathcal{E}, \quad (\mathcal{J}\Phi)(y) := \int_0^1 \Phi(sy)ds, \quad y \in E'; \\ \mathcal{I} : \mathcal{E}_{\log} &\rightarrow \mathcal{E}, \quad (\mathcal{I}\Phi)(y) := \int_0^1 \Phi(sy)s^{-1}ds, \quad y \in E'. \end{aligned} \quad (5)$$

Note that  $\mathcal{J}$  is well defined on all of  $\mathcal{E}$ , since by (3),  $\mathcal{J}\Phi$  is the Lévy exponent of the well-defined integral  $\int_{(0,1]} tdY(t)$ , where  $Y(1)$  has the Lévy exponent  $\Phi$ ; cf. Jurek (1985) or (2004). On the other hand,  $\mathcal{I}$  is only defined on  $\mathcal{E}_{\log}$ , which corresponds to infinitely divisible measures with finite

logarithmic moments, since  $\mathcal{I}\Phi$  is the Lévy exponent of the random integral  $\int_{(0,1]} t dY(-\ln t) = \int_{(0,\infty)} e^{-s} dY(s)$ , where  $\Phi$  is the Lévy exponent of  $Y(1)$  that has finite logarithmic moment; cf. Jurek-Vervaat (1983).

Here are the main algebraic properties of the mappings  $\mathcal{J}$  and  $\mathcal{I}$ .

**LEMMA 1.** *The operators  $\mathcal{I}$  and  $\mathcal{J}$  acting on appropriate domains (Lévy exponents) have the following basic properties:*

(a)  $\mathcal{I}, \mathcal{J}$  are additive and positive homogeneous operators;

(b)  $\mathcal{I}, \mathcal{J}$  commute under the composition and  $\mathcal{J}(\mathcal{I}(\Phi)) = (\mathcal{I} - \mathcal{J})\Phi$ .

*Other equivalent forms of that last property are:*

$$\mathcal{J}(I + \mathcal{I}) = \mathcal{I}; \quad \mathcal{I}(I - \mathcal{J}) = \mathcal{J}; \quad (I - \mathcal{J})(I + \mathcal{I}) = I.$$

*Proof.* Part (a) follows from the fact that  $\mathcal{E}$  forms a cone. For part (b) note that

$$\begin{aligned} (\mathcal{J}(\mathcal{I}(\Phi)))(y) &= \int_0^1 (\mathcal{I}(\Phi))(ty) dt = \int_0^1 \int_0^1 \Phi(sty) s^{-1} ds dt = \\ &= \int_0^1 \int_0^t \Phi(ry) r^{-1} dr dt = \int_0^1 \int_r^1 \Phi(ry) dt r^{-1} dr = \\ &= \int_0^1 \Phi(ry) r^{-1} dr - \int_0^1 \Phi(ry) dr = \mathcal{I}\Phi(y) - \mathcal{J}\Phi(y) = (\mathcal{I} - \mathcal{J})\Phi(y), \end{aligned}$$

which proves the equality in (b). Note that from the above (the first line of the above argument) we infer also that that operators  $\mathcal{I}$  and  $\mathcal{J}$  commute which completes the argument.  $\square$

**LEMMA 2.** *The operators  $\mathcal{I}$  and  $\mathcal{J}$ , defined by (5), have the following additional properties:*

(a)  $\mathcal{J} : \mathcal{E}_{\log} \rightarrow \mathcal{E}_{\log}$  and  $\mathcal{I} : \mathcal{E}_{(\log)^2} \rightarrow \mathcal{E}_{\log}$ ,

(b) *If  $(I - \mathcal{J})\Phi \in \mathcal{E}$  then the corresponding infinitely divisible measure  $\tilde{\mu}$  with the Lévy exponent  $(I - \mathcal{J})\Phi(y)$ ,  $y \in E'$ , has finite logarithmic moment.*

(c)  $(I - \mathcal{J})\Phi + \mathcal{I}(I - \mathcal{J})\Phi = (I - \mathcal{J})\Phi + \mathcal{J}\Phi = \Phi$  for all  $\Phi \in \mathcal{E}$ .

*Proof.* (a) Since the function  $E \ni x \rightarrow \log(1 + \|x\|)$  is sub-additive, for an infinitely divisible probability measure  $\mu = [a, R, M]$  we have

$$\begin{aligned} \int_E \log(1 + \|x\|) \mu(dx) < \infty & \text{ iff } \int_{\{\|x\| > 1\}} \log(1 + \|x\|) M(dx) < \infty \\ & \text{ iff } \int_{\{\|x\| > 1\}} \log \|x\| M(dx) < \infty; \end{aligned} \quad (6)$$

cf. Jurek and Mason (1993), Proposition 1.8.13. Furthermore, if  $M$  is the spectral Lévy measure appearing in the Lévy exponent  $\Phi$  then  $\mathcal{J}\Phi$  has Lévy spectral measure  $\mathcal{J}M$  (we keep that potentially conflicting notation), where

$$(\mathcal{J}M)(A) := \int_{(0,1)} M(t^{-1}A) dt = \int_{(0,1)} \int_E 1_A(tx) M(dx) dt, \quad (7)$$

for all Borel subsets  $A$  of  $E \setminus \{0\}$ . Hence

$$\begin{aligned} \int_{\{\|x\| > 1\}} \log \|x\| (\mathcal{J}M)(dx) &= \int_{(0,1)} \int_E 1_{\{\|x\| > 1\}}(tx) \log(t\|x\|) M(dx) dt \\ &= \int_{(0,1)} \int_{\{\|x\| > t^{-1}\}} \log(t\|x\|) M(dx) dt = \int_{\{\|x\| > 1\}} \int_{\|x\|^{-1}}^1 \log(t\|x\|) dt M(dx) \\ &= \int_{\{\|x\| > 1\}} \|x\|^{-1} \int_1^{\|x\|} \log w dw M(dx) \\ &= \int_{\{\|x\| > 1\}} \|x\|^{-1} [\|x\| \log \|x\| - \|x\| + 1] M(dx) \\ &= \int_{\{\|x\| > 1\}} \log \|x\| M(dx) - \int_{\{\|x\| > 1\}} [1 - \|x\|^{-1}] M(dx). \end{aligned}$$

Since the last integral is always finite as we integrate a bounded function with respect to a finite measure, we get the first part of (a). For the second one, let us note that

$$\begin{aligned} \int_{\{\|x\| > 1\}} \log \|x\| (\mathcal{I}M)(dx) &= \int_0^\infty \int_{\{\|x\| > 1\}} \log \|x\| M(e^t dx) dt \\ &= 1/2 \int_{\{\|x\| > 1\}} \log^2 \|x\| M(dx), \end{aligned}$$

where  $\mathcal{I}M$  is the Lévy spectral measure corresponding to the Lévy exponent  $\mathcal{I}\Phi$ .

For the part (b), note that the assumption made there implies that the measure

$$\widetilde{M}(A) := M(A) - \int_{(0,1)} M(t^{-1}A)dt \geq 0, \text{ for all Borel sets } A \subset E \setminus \{0\}, \quad (8)$$

is the Lévy spectral measure of some  $\tilde{\mu}$ . [Note that there is no restriction on the Gaussian part.] In fact, if  $\widetilde{M}$  is a nonnegative measure then it is necessarily a Lévy spectral measure because  $0 \leq \widetilde{M} \leq M$  and  $M$  is Lévy spectral measure; comp. Arujo-Giné (1980), Chapter 3, Theorem 4.7 , p. 119. To establish the logarithmic moment of  $\tilde{\mu}$  we argue as follows. Observe

that for any constant  $k > 1$  we have

$$\begin{aligned}
0 &\leq \int_{\{1 < \|x\| \leq k\}} \log \|x\| \widetilde{M}(dx) = \\
&\int_{\{1 < \|x\| \leq k\}} \log \|x\| M(dx) - \int_{(0,1)} \int_{\{1 < \|x\| \leq k\}} \log \|x\| M(t^{-1}dx) dt = \\
&\int_{\{1 < \|x\| \leq k\}} \log \|x\| M(dx) - \int_{(0,1)} \int_{\{t^{-1} < \|x\| \leq kt^{-1}\}} \log(t\|x\|) dM(dx) dt = \\
&\int_{\{1 < \|x\| \leq k\}} \log \|x\| M(dx) - \int_{\{1 < \|x\| \leq k\}} \int_{\|x\|^{-1}}^1 \log(t\|x\|) dt M(dx) \\
&\quad - \int_{\{k < \|x\|\}} \int_{\|x\|^{-1}}^{k\|x\|^{-1}} \log(t\|x\|) dt M(dx) = \\
&\int_{\{1 < \|x\| \leq k\}} \log \|x\| M(dx) - \int_{\{1 < \|x\| \leq k\}} \|x\|^{-1} \int_1^{\|x\|} \log(w) dw M(dx) \\
&\quad - \int_{\{k < \|x\|\}} \|x\|^{-1} \int_1^k \log(w) dw M(dx) = \\
&\int_{\{1 < \|x\| \leq k\}} \log \|x\| M(dx) - \int_{\{1 < \|x\| \leq k\}} \|x\|^{-1} (\|x\| \log \|x\| - \|x\| + 1) M(dx) \\
&\quad - (k \log k - k + 1) \int_{\{\|x\| > k\}} \|x\|^{-1} M(dx) = \\
&\int_{\{1 < \|x\| \leq k\}} (1 - \|x\|^{-1}) M(dx) - (k \log k - k + 1) \int_{\{\|x\| > k\}} \|x\|^{-1} M(dx) \\
&\leq M(\|x\| > 1) < \infty,
\end{aligned}$$

and consequently  $\int_{(\|x\| > 1)} \log \|x\| \widetilde{M}(dx) < \infty$ . This with property (6), completes the proof of the part (b).

Finally, since  $(I - \mathcal{J})\Phi$  is in a domain of definition of the operator  $\mathcal{I}$ , so the part (c) is a consequence of Lemma 1(e) and (d). Thus the proof is complete.  $\square$

**3. Factorizations of selfdecomposable distributions.** The classes of limit laws  $\mathcal{U}$  and  $L$  are obtained by non-linear shrinking transformations and linear transformations (multiplications by scalars), respectively; cf. Jurek (1985) and references there. However, there are many (unexpected) relations

between  $\mathcal{U}$  and  $L$  as was already proved in Jurek (1985) and more recently in Iksanov-Jurek-Schreiber (2004). Furthermore, more recently selfdecomposable distributions are used in modelling real phenomena, in particular in mathematical finance; for instance cf. Bingham (2006), Carr-Geman-Madan-Yor (2005) or Eberlein-Keller (1995). This motivates further studies on factorizations and other relations between the classes  $\mathcal{U}$  and  $L$ , like those in Theorems 1 and 2, below.

In this section we will apply the operators  $\mathcal{I}$  and  $\mathcal{J}$  to Lévy exponents of selfdecomposable (the class  $L$ ) and s-selfdecomposable (the class  $\mathcal{U}$ ) probability measures. For the convenience of the readers recall here that

$$\begin{aligned} \mu \in L \quad &\text{iff} \quad \forall(t > 0) \exists \nu_t \quad \mu = T_{e^{-t}} \mu * \nu_t \\ &\text{iff} \quad \mu = \mathcal{L}\left(\int_{(0,\infty)} e^{-t} dY(t)\right); \quad \mathcal{L}(Y(1)) \in ID_{\log}, \\ \mu \in \mathcal{U} \quad &\text{iff} \quad \mu = \mathcal{L}\left(\int_{(0,1]} t dY(t)\right), \quad \mathcal{L}(Y(1)) \in ID. \end{aligned} \tag{9}$$

Measures from the class  $\mathcal{U}$  are called *s-selfdecomposable*; cf Jurek (1985), (2004). The corresponding Fourier transforms of measures from  $L$  and  $\mathcal{U}$  easily follow from (2) and (3); cf. Jurek-Vervaat (1983) or the above references.

**LEMMA 3.** *If  $\mu$  is a selfdecomposable probability measure on a Banach space  $E$  with characteristic function  $\hat{\mu}(y) = \exp[\Phi(y)]$   $y \in E'$ , then*

$$\tilde{\Phi}(y) := \Phi(y) - \int_{(0,1)} \Phi(sy) ds = (I - \mathcal{J})\Phi(y), \quad y \in E',$$

*is a Lévy exponent corresponding to an infinitely divisible probability measure with finite logarithmic moment.*

*Equivalently, if  $\widetilde{M}$  is the Lévy spectral measure of a selfdecomposable  $\mu$  then the measure  $\widetilde{M}$  given by*

$$\widetilde{M}(A) := M(A) - \int_0^1 M(t^{-1}A) dt, \quad A \subset E \setminus \{0\},$$

*is a Lévy spectral measure on  $E$  that additionally integrates the logarithmic function on the complement of any neighborhood of zero.*



*Proof.* If  $\mu = [a, R, M]$  is selfdecomposable (or in other words a class L distribution) then we infer that

$$M(A) - M(e^t A) \geq 0, \text{ for all } t > 0 \text{ and Borel } A \subset E \setminus \{0\},$$

and that there is no restriction on the remaining two parameters (the shift vector and the Gaussian covariance operator) in the Lévy-Khintchine formula (1). Multiplying both sides by  $e^{-t}$  and then integrating over the positive half-line we conclude that  $\widetilde{M}$ , given by (8), is a non-negative measure. Since  $\widetilde{M} \leq M$  and  $M$  is a Lévy spectral measure, so is  $\widetilde{M}$ ; comp. Theorem 4.7 in Chapter 3 of Araujo-Giné (1980). Finally, our Lemma 2(b) gives the finiteness of the logarithmic moment. Thus the proof is complete.  $\square$

**THEOREM 1.** *For each selfdecomposable probability measure  $\mu$ , on a Banach space  $E$ , there exists a unique  $s$ -selfdecomposable probability measure  $\tilde{\mu}$  with finite logarithmic moment such that*

$$\mu = \tilde{\mu} * \mathcal{I}(\tilde{\mu}) \text{ and } \mathcal{J}(\mu) = \mathcal{I}(\tilde{\mu}). \quad (10)$$

*In fact, if  $\hat{\mu}(y) = \exp[\Phi(y)]$  then  $(\tilde{\mu})(y) = \exp[\Phi(y) - \int_{(0,1)} \Phi(ty)dt]$ ,  $y \in E'$ .*

*In other words, if  $\Phi$  is the Lévy exponent of a selfdecomposable probability measure then  $(I - \mathcal{J})\Phi$  is the Lévy exponent of an  $s$ -selfdecomposable measure with the finite logarithmic moment and*

$$\Phi = (I - \mathcal{J})\Phi + \mathcal{I}(I - \mathcal{J})\Phi = (I - \mathcal{J})\Phi + \mathcal{J}\Phi. \quad (11)$$

*Proof.* Let  $\hat{\mu}(y) = \exp[\Phi(y)] \in L$ . From the factorization in (9) (the first line) we infer that  $\Phi_t(y) := \Phi(y) - \Phi(e^{-t}y)$  are Lévy exponents as well. Hence,

$$\tilde{\Phi}(y) := \int_{(0,\infty)} \Phi_t(ty)e^{-t}dt = \Phi(y) - \int_{(0,\infty)} \Phi(e^{-t}y)e^{-t}dt = ((I - \mathcal{J})\Phi)(y)$$

is a Lévy exponent as well, because of Lemma 3. Again by Lemma 3 (or Lemma 2 b)), a probability measure  $\tilde{\mu}$  defined by the Fourier transform  $(\tilde{\mu})(y) = \exp((I - \mathcal{J})\Phi(y))$  has logarithmic moment. Consequently,  $\mathcal{I}(\tilde{\mu})$  is a well defined probability measure whose Lévy exponent is equal to  $\mathcal{I}(I - \mathcal{J})\Phi$ . Finally, Lemmas 1(b) and 2(c) give the factorization (10).

Since  $\mathcal{I}(\tilde{\mu}) \in L$  has the property that  $\tilde{\mu} * \mathcal{I}(\tilde{\mu})$  is again in  $L$ , therefore Theorem 1 from Iksanov-Jurek-Schreiber(2004) gives that  $\tilde{\mu} \in \mathcal{U}$ , i.e., it is a  $s$ -selfdecomposable probability distribution.

To see the second equality in (11) one should observe that it is equivalent to equality  $\mathcal{J}\Phi = \mathcal{I}(I - \mathcal{J})\Phi$  that indeed holds true in view of Lemma 1(d).

Suppose there exists another factorization of the form  $\mu = \rho * \mathcal{I}(\rho)$  and let  $\Xi(y)$  be the Lévy exponent of  $\rho$ . Then we get that  $\Phi(y) = \Xi(y) + (\mathcal{I}\Xi)(y) = (I + \mathcal{I})\Xi(y)$ . Hence, applying to both sides  $\mathcal{I} - \mathcal{J}$  we conclude that

$$(I - \mathcal{J})\Phi = ((I - \mathcal{J})(I + \mathcal{I}))\Xi = \Xi,$$

where the last equality is from Lemma 1(b). This proves the uniqueness of  $\tilde{\mu}$  in the representation (10) and thus the proof of Theorem 1 is completed.  $\square$

*REMARK 2.* The factorization (10), in Theorem 1, can be also derived from previous papers as follows:

*for each selfdecomposable (or class L)  $\mu$  there exists a unique  $\rho \in ID_{\log}$  such that  $\mu = \mathcal{I}(\rho)$ ; Jurek-Vervaat (1983). Since  $\tilde{\mu} := \mathcal{J}(\rho)$  is s-selfdecomposable (class  $\mathcal{U}$ ) with logarithmic moment (cf. Jurek (1983)) therefore,  $\mathcal{I}(\tilde{\mu}) * \tilde{\mu} \in L$  in view of Iksanov-Jurek-Schreiber (2004). Finally, again by Jurek (1985),  $\mathcal{I}(\tilde{\mu}) * \tilde{\mu} = \mathcal{J}(\mathcal{I}(\rho) * \rho) = \mathcal{I}(\rho) = \mu$ , which gives the decomposition.*

However, the present proof is less involved, more straightforward and moreover the result and the proof of finiteness of the logarithmic moment in Lemma 2 (b) is completely new. Last but not least, the "calculus" on Lévy exponents, introduced in this note, is of an interest in itself.

*REMARK 3.* In the case of Euclidean space  $\mathbb{R}^d$ , using Schoenberg's Theorem, one gets immediately that  $\tilde{\Phi}$  is a Lévy exponent; cf. Cuppens (1975), pp. 80-82.

Following Iksanov, Jurek and Schreiber (2004), p. 1360, we will say that a selfdecomposable probability measure  $\mu$  has *the factorization property* if  $\mu * \mathcal{I}^{-1}(\mu)$  is selfdecomposable as well. In other words, a class  $L$  probability measure convolved with its background driving probability distribution is again class  $L$  distribution. As in Iksanov-Jurek-Schreiber (2004), Proposition 1, if  $L^f$  denotes the set of all class  $L$  distribution with the factorization property then

$$L^f = \mathcal{I}(\mathcal{J}(ID_{\log})) = \mathcal{J}(\mathcal{I}(ID_{\log})) = \mathcal{J}(L) \text{ and } L^f \subset L \subset \mathcal{U}, \quad (12)$$

**COROLLARY 1.** *Each selfdecomposable  $\mu$  admits a factorization  $\mu = \nu_1 * \nu_2$ , where  $\nu_1$  is an s-selfdecomposable measure (i.e.,  $\nu_1 \in \mathcal{U}$ ) and  $\nu_2$  is a selfdecomposable one with the factorization property (i.e.,  $\nu_2 \in L^f$ ). That is, besides the inclusion  $L^f \subset L \subset \mathcal{U}$  we also have that  $L \subset L^f * \mathcal{U}$ .*

*Proof.* Because of (10),  $\nu_1 := \tilde{\mu}$  is an s-selfdecomposable measure. Furthermore,  $\nu_2 := \mathcal{I}(\tilde{\mu}) \in L$  has the factorization property, i.e.,  $\nu_2 \in L^f$ , which completes the proof.  $\square$

**EXAMPLES.** 1) Let  $\Sigma_p$  be a symmetric stable distribution on a Banach space  $E$ , with the exponent  $p$ . Then its Lévy exponent,  $\Phi_p$ , is equal to  $\Phi_p(y) = -\int_S | \langle y, x \rangle |^p m(dx)$ , where  $m$  is a finite Borel measure on the unit sphere  $S$  of  $E$ ; cf. Samorodnitsky and Taqqu (1994). Hence  $(I - \mathcal{J})\Phi_p(y) = p/(p+1)\Phi_p(y)$ , which means that in Corollary 1, both  $\nu_1$  and  $\nu_2$  are stable with the exponent  $p$  and measures  $m_1 := (p/(p+1))m$  and  $m_2 := (1/(p+1))m$ , respectively.

2) Let  $\eta$  denotes the Laplace (double exponential) distribution on real line  $\mathbb{R}$ ; cf. Jurek-Yor (2004). Then its Lévy exponent  $\Phi_\eta$  is equal to  $\Phi_\eta(t) := -\log(1+t^2)$ ,  $t \in \mathbb{R}$ . Consequently,  $(I - \mathcal{J})\Phi_\eta(t) = 2(\arctan t - t)t^{-1}$  is the Lévy exponent of the class  $\mathcal{U}$  probability measure  $\nu_1$  from Corollary 1, and  $(2t - \arctan t - t \log(1+t^2))t^{-1}$  is the Lévy exponent of the class  $L^f$  measure  $\nu_2$  from Corollary 1.

Before we formulate the next result we need to recall that, by (9), the class  $\mathcal{U}$  is defined here as  $\mathcal{U} = \mathcal{J}(ID)$ . Consequently, by iteration argument we can define

$$\mathcal{U}^{<1>} := \mathcal{U}, \quad \mathcal{U}^{<k+1>} := \mathcal{J}(\mathcal{U}^{<k>}) = \mathcal{J}^{k+1}(ID), \quad k = 1, 2, \dots; \quad (13)$$

cf. Jurek (2004) for other characterization of classes  $\mathcal{U}^{<k>}$ . Elements from the semigroup  $\mathcal{U}^{<k>}$  are called *k-times s-selfdecomposable probability measures*.

**THEOREM 2.** *Let  $n$  be any natural number and  $\mu$  be a selfdecomposable probability measure. Then there exist k-times s-selfdecomposable probability measures  $\tilde{\mu}_k$ , for  $k = 1, 2, \dots, n$ , such that*

$$\mu = \tilde{\mu}_1 * \tilde{\mu}_2 * \dots * \tilde{\mu}_n * \mathcal{I}(\tilde{\mu}_n), \quad \mathcal{J}^k(\mu) = \mathcal{I}(\tilde{\mu}_k), \quad k = 1, 2, \dots, n. \quad (14)$$

*In fact, if  $\Phi$  is the exponent of  $\mu$  then  $\tilde{\mu}_k$  has the exponent  $\mathcal{I}^{k-1}(I - \mathcal{J})^k\Phi = (I - \mathcal{J})\mathcal{J}^{k-1}\Phi$  and*

$$\begin{aligned} \Phi &= (I - \mathcal{J})\Phi + (I - \mathcal{J})\mathcal{J}\Phi + \dots + (I - \mathcal{J})\mathcal{J}^{k-1}\Phi + \dots + (I - \mathcal{J})\mathcal{J}^{n-1}\Phi + \mathcal{J}^n\Phi \\ &= (I - \mathcal{J}^n)\Phi + \mathcal{J}^n\Phi. \end{aligned} \quad (15)$$

*Proof.* For  $n = 1$  the factorization (14) and the formula (15) are true by Theorem 1, with  $\tilde{\mu}_1 := \tilde{\mu}$ . Suppose our claim (14) is true for  $n$ . Since  $\rho := \mathcal{I}(\tilde{\mu}_n)$  is selfdecomposable, applying to it Theorem 1, we have that  $\rho = \tilde{\rho} * \mathcal{I}(\tilde{\rho})$ , where  $\tilde{\rho}$  has the Lévy exponent  $(I - \mathcal{J})\mathcal{J}^n\Phi = \mathcal{J}^n(I - \mathcal{J})\Phi$  and thus it corresponds to  $(n + 1)$ -times s-selfdecomposable probability because, by Theorem 1,  $(I - \mathcal{J})\Phi$  is already s-selfdecomposable and then we apply  $n$  times the operator  $\mathcal{J}$ ; compare the definition (13). Thus the factorization (14) holds for  $n + 1$ , which completes the proof of the first part of the theorem. Similarly, applying inductively decomposition (11), from Theorem 1 and observing from Lemma 1(b) that we will get the formula (14). Thus the proof is complete.  $\square$

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